Tropical heat The eikonal equation as a (max,+) version of the Poisson equation

Averil Prost

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Finite-dim systems

Table of Contents

The (max,+) semialgebra

Link with linear algebra

An application: finite-dimensional (max, +) system

The heat equation

Finite-dim systems

The (max,+) idempotent calculus

Let $\Omega \coloneqq \mathbb{R} \cup \{-\infty\}$.

Def 1 – Operations For $a, b \in \Omega$, define

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Finite-dim systems

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Finite-dim systems

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Finite-dim systems

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"Idempotent" since $a \oplus a = a$. Define $\mathbb{O} := -\infty$ and $\mathbb{I} := 0$. Then

$$\mathbb{0} \oplus a = \max\left(-\infty, a\right) = a, \qquad \mathbb{0} \otimes a = -\infty + a = \mathbb{0}, \qquad \mathbb{1} \otimes a = a + 0 = a.$$

Then $(\Omega, \oplus, \otimes, \mathbb{0}, \mathbb{1})$ is a semiring (ring without additive inverse).

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The (max,+) semialgebra ००●०	Link with linear algebra	Finite-dim systems 000000	The heat equation
Integrals			

Endow Ω with the application $d(a,b) \coloneqq |e^a - e^b|$, and let $f: \Omega \to \Omega$ be continuous. Then

$$\sum_{i\in\mathbb{Z}\cup\{-\infty\}}^{\oplus} f(hi) = \max_{i\in\mathbb{Z}\cup\{-\infty\}} f(hi) \xrightarrow[h\searrow 0]{} \sup_{x\in\Omega=\mathbb{R}\cup\{-\infty\}} f(x).$$

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Def 2 – Integral Define $\int_{x\in\Omega}^\oplus f(x)\coloneqq \sup_{x\in\Omega} f(x).$

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Def 2 – Integral Define
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In particular, the scalar product becomes

$$\langle f,g \rangle_{\oplus} = \int_{x \in \Omega}^{\oplus} f(x) \otimes g(x) = \sup_{x \in \Omega} f(x) + g(x).$$

Finite-dim systems

People and motivation

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Motivation Ease the study of **optimization problems** by directly working with the "natural" operations of (here) **maximization and sum** of gains.

Finite-dim systems

The heat equation

Table of Contents

The (max,+) semialgebra

Link with linear algebra

An application: finite-dimensional (max, +) system

The heat equation

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	o●oooo	000000	

An elementary result

Proposition –
 Let
$$f > g$$
. Then

$$\lim_{h \searrow 0} h \log \left(e^{f/h} + e^{g/h} \right) = \lim_{h \searrow 0} h \log \left(e^{f/h} - e^{g/h} \right) = f.$$

The (max,+) semialgebra 0000	Link with linear algebra o●oooo	Finite-dim systems	The heat equation

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Readily visible by noticing that

$$e^{f/h} \pm e^{g/h} = e^{f/h} \left(1 \pm e^{(g-f)/h} \right).$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	o●oooo	000000	

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Going further, for any upper bounded and continuous f, there holds

$$\lim_{h\searrow 0} h\log\left(\int_{x\in \mathbb{R}^d} \exp\left(\frac{f(x)}{h}\right) dx\right) = \sup_{x\in \mathbb{R}^d} f(x).$$

The ((max,+)	semialgebra

Finite-dim systems

The heat equation

Logarithm transform

Def 3 – **Logarithm trick** For h > 0, consider the operations

$$a \otimes_h b \coloneqq h \log \left(e^{a/h} \cdot e^{b/h} \right), \qquad a \oplus_h b \coloneqq h \log(e^{a/h} + e^{b/h}).$$

	The (max,+) semialgebra	Link with linear algebra ००●०००	Finite-dim systems 000000	
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Then

$$a \otimes_h b = a + b, \qquad a \oplus_h b \xrightarrow[h \searrow 0]{} \max(a, b).$$

The (max,+) semialgebra	Link with linear algebra 00●000	Finite-dim systems 000000	The heat equatic

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We could go further with this game, for instance with

$$a^{\oplus}b \coloneqq \lim_{h \searrow 0} h \log \left(\exp(a/h) \circ \exp(b/h) \right)$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The he
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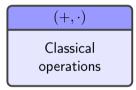
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$$a^{\oplus}b \coloneqq \lim_{h \searrow 0} h \log \left(\exp(a/h) \circ \exp(b/h) \right) = \begin{cases} a & b = 0\\ \mathbb{I} = 0 & b < 0\\ +\infty & b > 0. \end{cases}$$

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Finite-dim systems

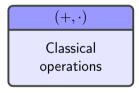
The heat equation



Finite-dim systems

The heat equation 000000000

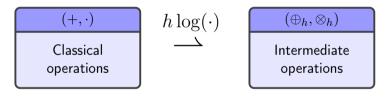
Dialog between $(+,\cdot)$ and $(\max,+)$



 (\oplus_h,\otimes_h) Intermediate operations

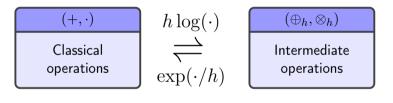
Finite-dim systems

The heat equation 000000000



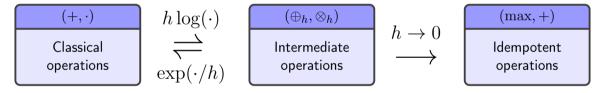
Finite-dim systems

The heat equation



Finite-dim systems

The heat equation



The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Linearity

Let $\mathrm{I}:\mathbb{R}^d \to \mathbb{R}^d$ be linear, and denote

 $\Phi_h(x) \coloneqq h \log \circ \mathbf{I} \circ \exp\left(x/h\right).$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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The (max,+) semialgebra 0000	Link with linear algebra ००००●੦	Finite-dim systems	The heat equation

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Similarly,

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The (max,+) semialgebra	Link with linear algebra ୦୦୦୦●୦	Finite-dim systems 000000	The heat equation

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The limit operator $\Phi \coloneqq \lim_{h \searrow 0} \Phi_h$ is $(\max, +)$ -linear.

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Example			

Let
$$d = 2$$
 and $I(x) = Ax$, where $A \coloneqq \begin{pmatrix} 1 & 1 \\ 2 & 0 \end{pmatrix}$. Then

 $\Phi_h(x) = h \log \circ \mathbf{I} \circ \exp\left(x/h\right)$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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The (max,+) semialgebra	Link with linear algebra 00000●	Finite-dim systems 000000	The heat equation
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$\Phi_h(x) = h \log \circ \mathbf{I} \circ \exp\left(x\right)$	$(x/h) = h \log \left(\exp(x_1/h) + 2\exp(x_2/h) \right)$	$\left(\begin{array}{c} \exp\left(x_2/h\right) \\ x_1/h \end{array} \right) \xrightarrow[h \searrow 0]{} \left(\begin{array}{c} x_1 \\ x_2 \\ x_1 \end{array} \right)$	$\left. \begin{array}{c} \max(x_1, x_2) \\ x_1 \end{array} \right).$

The (max,+) semialgebra 0000	Link with linear algebra 00000●	Finite-dim systems 000000	The heat equation
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$\Phi_h(x) = h \log \circ \mathbf{I} \circ \exp\left(x\right)$	$(x/h) = h \log \left(\exp \left(\frac{x_1}{2e} \right) \right)$	$(h) + \exp(x_2/h) \xrightarrow{h \searrow 0} \left(\exp(x_1/h) \right) \xrightarrow{h \searrow 0} \left(\left(\frac{h}{h} \right) \right)$	$\begin{pmatrix} \max(x_1, x_2) \\ x_1 \end{pmatrix}.$
Define $\Phi(x) \coloneqq \begin{pmatrix} \max(x_1, x_1) \\ x_1 \end{pmatrix}$	$\left(\mathcal{S}_{2} ight) ight) :$ then for any $\lambda \in$	$\mathbb R$, (x_1,x_2) and (y_1,y_2) ,	

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$\Phi\left(\lambda\otimes x\oplus y\right)=\Phi\left(\begin{matrix}\max_{\max}\\\\\max\end{matrix}\right)$	$egin{aligned} & (\lambda+x_1,y_1) \ & (\lambda+x_2,y_2) \end{aligned} ight)$		

The (max,+) semialgebra 0000	Link with linear algebra 00000●	Finite-dim systems 000000	The heat equation
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$\Phi_h(x) = h \log \circ \mathbf{I} \circ \exp \left(-\frac{1}{2} \right)$	$(x/h) = h \log \left(\exp \left(x_1 \right) \right)$	$(h) + \exp(x_2/h) $ $\xrightarrow{h \searrow 0} $ (1)	$\left. \begin{array}{c} \max(x_1, x_2) \\ x_1 \end{array} \right).$
Define $\Phi(x) \coloneqq \begin{pmatrix} \max(x_1, \\ x_1 \end{pmatrix}$	$x_2) \Biggr)$: then for any $\lambda \in$	\mathbb{R} , (x_1,x_2) and (y_1,y_2) ,	
$\Phi\left(\lambda\otimes x\oplus y ight)=\Phi\left(egin{max}{max}(max)\max ight)$	$\begin{pmatrix} \lambda + x_1, y_1 \end{pmatrix} \\ \lambda + x_2, y_2 \end{pmatrix} = \begin{pmatrix} \max \\ \end{pmatrix}$	$ \begin{pmatrix} \lambda + x_1, y_1, \lambda + x_2, y_2 \end{pmatrix} \\ \max(\lambda + x_1, y_1) \end{pmatrix} $	

The (max,+) semialgebra	Link with linear algebra 00000●	Finite-dim systems	The heat equation
Example			
Let $d = 2$ and $I(x) = Ax$, we have:	where $A\coloneqq egin{pmatrix} 1 & 1 \ 2 & 0 \end{pmatrix}$. Th	en	
$\Phi_h(x) = h \log \circ \mathbf{I} \circ \exp\left(x\right)$	$(x/h) = h \log \left(\exp \left(\frac{x_1/h}{2 \exp \left(\frac{x_2}{2} + \frac{x_1}{2} + \frac{x_1}{2} + \frac{x_1}{2} + \frac{x_2}{2} + \frac{x_1}{2} + \frac{x_2}{2} + \frac{x_1}{2} + \frac{x_2}{2} + \frac{x_1}{2} + x_$	$(h) + \exp(x_2/h) \to (x_1/h) \to (h \searrow 0)$	$\begin{pmatrix} \max(x_1, x_2) \\ x_1 \end{pmatrix}.$
Define $\Phi(x) \coloneqq \begin{pmatrix} \max(x_1, x_1) \\ x_1 \end{pmatrix}$	$\left(x_{2} ight) ight)$: then for any $\lambda\in\mathbb{R}$	(x_1,x_2) and (y_1,y_2) ,	
$\Phi\left(\lambda\otimes x\oplus y\right)=\Phi\left(\max_{\max(\lambda)}\right)$	$\begin{pmatrix} \lambda + x_1, y_1 \end{pmatrix} = \begin{pmatrix} \max(\lambda \\ n \end{pmatrix}$	$\left(\lambda + x_1, y_1, \lambda + x_2, y_2\right)$ $\max(\lambda + x_1, y_1)$	$=\lambda \otimes \Phi(x) \oplus \Phi(y).$

Averil Prost

Link with linear algebra

Finite-dim systems

The heat equation

Table of Contents

The (max,+) semialgebra

Link with linear algebra

An application: finite-dimensional $(\max,+)$ system

The heat equation

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems 0●0000

The heat equation

Setting

Consider the (max, +) system

$$\Pi \otimes \xi = \beta, \tag{1}$$

where $\Pi \in \mathbb{M}_{2,2}$ is a matrix, $\xi, \beta \in \mathbb{R}^2$ are vectors with β given, and for each $i \in \llbracket 1, 2 \rrbracket$,

$$(\Pi \otimes \xi)_i = \sum_{j \in \llbracket 1, 2 \rrbracket}^{\oplus} \Pi_{ij} \otimes \xi_j = \max_{j \in \llbracket 1, 2 \rrbracket} \Pi_{ij} + \xi_j.$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems 0●0000	The heat equation

Setting

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In which cases can we get a solution to (1) by using the link with classical algebra?

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Transformation into linear variables

For each h > 0, define $P \in \mathbb{M}_{2,2}$ and $x, b \in \mathbb{R}^2$ by

$$P_{ij}^h \coloneqq \exp(\Pi_{ij}/h), \qquad x_j^h \coloneqq \exp(\xi_j/h), \qquad b_i^h \coloneqq \exp(\beta_i/h).$$

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The linear system $P^h x^h = b^h$ has a solution if $\det(P^h) \neq 0,$ i.e.

$$P_{11}^h P_{22}^h - P_{12}^h P_{21}^h \neq 0 \qquad \iff \qquad \exp\left(\frac{\Pi_{11} + \Pi_{22}}{h}\right) \neq \exp\left(\frac{\Pi_{12} + \Pi_{21}}{h}\right).$$

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems 000●00	The heat equation
Going back			

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Going back			

$$\forall i \in [\![1,2]\!], \quad x_i^h = \left((P^h)^{-1} b \right)_i \ge 0.$$
 (2)

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Going back			

$$\forall i \in [\![1,2]\!], \quad x_i^h = \left((P^h)^{-1} b \right)_i \ge 0.$$
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One also obtains that $(x^h)_h$ is bounded uniformly in h. By compactness, (1) admits solutions.

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	000●00	
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$$\Pi = \begin{pmatrix} \mathbb{I} & \mathbb{0} \\ \mathbb{0} & \mathbb{I} \end{pmatrix} = \begin{pmatrix} 0 & -\infty \\ -\infty & 0 \end{pmatrix}$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Going back			

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	000●00	
Going back			

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Then $P^{-1} = \mathbb{I}_d$. As (2) is satisfied for all $b = \exp(\beta/h)$, we obtain that $\xi = \beta$ solves (1).

The ((max,+)	semialg	ebra

Link with linear algebra

Finite-dim systems 0000●0

The heat equation

Another example

${\sf Consider}$

$$\Pi = \begin{pmatrix} 1 & 2 \\ 4 & 3 \end{pmatrix}$$

The ((max,+)	semialgebra

Link with linear algebra

Finite-dim systems 0000●0

The heat equation

Another example

Consider

$$\Pi = \begin{pmatrix} 1 & 2 \\ 4 & 3 \end{pmatrix} \implies P = \begin{pmatrix} e^{1/h} & e^{2/h} \\ e^{4/h} & e^{3/h} \end{pmatrix},$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	0000●0	

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Since $\det P^h = e^{4/h} - e^{6/h} < 0$, the condition (2) over β becomes

$$e^{\frac{3+\beta_1}{h}}-e^{\frac{4+\beta_2}{h}}\leqslant 0, \quad \text{and} \quad -e^{\frac{2+\beta_1}{h}}+e^{\frac{1+\beta_2}{h}}\leqslant 0.$$

The (ma x,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	0000●0	

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0000	000000	0000●0	

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Whenever this condition is satisfied, the solution is given by

$$\begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \lim_{h \searrow 0} \begin{pmatrix} h \log \left(e^{\frac{4+\beta_2}{h}} - e^{\frac{3+\beta_1}{h}} \right) - h \log \left(e^{6/h} - e^{4/h} \right) \\ h \log \left(e^{\frac{2+\beta_1}{h}} - e^{\frac{1+\beta_2}{h}} \right) - h \log \left(e^{6/h} - e^{4/h} \right) \end{pmatrix}$$

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	0000●0	

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The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
0000	000000	0000●0	

Another example

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The ((max,+)	semialgebra	

Link with linear algebra

Finite-dim systems 00000●

The heat equation

To go further

Of course, the bulky reasoning may be refined.

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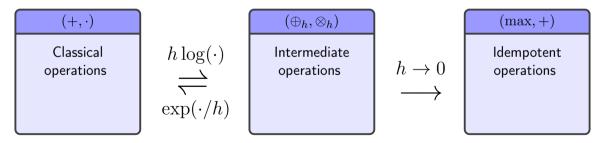
• Algorithms to solve linear problems may be transposed in the (max, +) semialgebra (for instance, Euler scheme ⇒ the semi-lagrangian). Even tropical finite elements [McE06]!

Finite-dim systems

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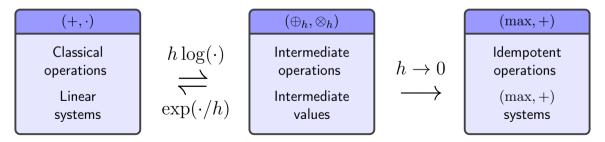
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Finite-dim systems

Table of Contents

The (max,+) semialgebra

Link with linear algebra

An application: finite-dimensional $(\max, +)$ system

The heat equation

 $\ensuremath{\operatorname{FIRST}}$ ORDER Consider the first-order transport equation

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Then

$$\partial_t u = \frac{\exp(v/h)}{h} \partial_t v, \quad \nabla u = \frac{\exp(v/h)}{h} \nabla v,$$
(3)

 $\label{eq:First-order} First \text{ order } the \ first \text{-order } transport \ equation$

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so that after dividing by $\frac{\exp(v/h)}{h} > 0$, we get again $\partial_t v(t,x) - \langle \nabla v(t,x), b(t,x) \rangle = 0$.

The (max,+) semialgebra 0000	Link with linear algebra	Finite-dim systems	The heat equation 00●000000

 Second order $% \operatorname{Consider}$ how the heat equation

 $\partial_t u(t,x) - h\Delta u(t,x) = 0.$

The (max,+) semialgebra	Link with linear algebra 000000	Finite-dim systems	The heat equation

 $\begin{array}{c|c} {\rm Second \ order} \end{array} \mbox{ Consider now the heat equation} \\ \end{array}$

$$\partial_t u(t,x) - h\Delta u(t,x) = 0.$$

Let again $v = h \log(u)$, i.e. $u = \exp(v/h)$. Then in addition to (3), there holds

$$\Delta u = \frac{\exp(v/t)}{h} \Delta v + \frac{\exp(v/t)}{h^2} |\nabla v|^2.$$

The (max,+) semialgebra 0000	Link with linear algebra 000000	Finite-dim systems	The heat equation

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Hence

$$\frac{\exp(v/h)}{h}\partial_t v - h\left(\frac{\exp(v/t)}{h}\Delta v + \frac{\exp(v/t)}{h^2}\left|\nabla v\right|^2\right) = 0,$$

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Hence

$$\frac{\exp(v/h)}{h}\partial_t v - h\left(\frac{\exp(v/t)}{h}\Delta v + \frac{\exp(v/t)}{h^2}\left|\nabla v\right|^2\right) = 0,$$

or after simplification,

$$\partial_t v(t,x) - h\Delta v(t,x) - |\nabla v(t,x)|^2 = 0.$$

The (max,+) semialgebra	Link with linear algebra 000000	Finite-dim systems 000000	The heat equation

A link with viscosity solutions

Def 4 – (Historical intuition, see [CL83]) The vanishing viscosity solution of $\partial_t v - |\nabla v|^2 = 0$

is the limit when h goes to 0 of the (unique) solution of the equation

 $\partial_t v - h\Delta v - |\nabla v|^2 = 0.$

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A link with viscosity solutions

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 $\partial_t v - |\nabla v|^2 = 0$

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It may be characterized by two sign inequalities that maintain the validity of the comparison principle coming from the elliptic perturbation.

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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The heat kernel

Def 5 – Heat kernel Let C_d be a normalizing constant, and for any $(t, x) \in \mathbb{R}^+ \times \mathbb{R}^d$, define

$$\mu_{t,x}^h \coloneqq \frac{1}{\sqrt{S_d t h}} \exp\left(-\frac{|\cdot - x|^2}{2th}\right) \mathcal{L}_{\mathbb{R}^d} \in \mathscr{P}_2(\mathbb{R}^d).$$

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In particular, $\mu^h_{t,x} \underset{t\searrow 0}{\longrightarrow} \delta_x$ narrowly and in the Wasserstein topology.

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In particular, $\mu^h_{t,x} \underset{t\searrow 0}{\longrightarrow} \delta_x$ narrowly and in the Wasserstein topology.

Proposition – Kernel representation The (weak) solution of the heat equation $\partial_t u(t,x) - h\Delta u(t,x) = 0$ with initial value u_0 is given by

$$u(t,x) = \int_{y \in \mathbb{R}^d} u_0(y) d\mu^h_{t,x}(y).$$

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Finite-dim systems

Going through the transformation

Let $u_0 > 0$ upper bounded and continuous.

Finite-dim systems

The heat equation

Going through the transformation

Let $u_0 > 0$ upper bounded and continuous. Denoting again $v = h \log(u)$, and $v_0 = h \log(u_0)$,

$$v(t,x) = h \log \left(\int_{y \in \mathbb{R}^d} \exp \left(\frac{v_0(y)}{h} \right) d\mu_{t,x}^h(y) \right)$$

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$$\begin{aligned} v(t,x) &= h \log \left(\int_{y \in \mathbb{R}^d} \exp\left(\frac{v_0(y)}{h}\right) d\mu_{t,x}^h(y) \right) \\ &= h \log\left(\frac{1}{\sqrt{C_d t h}} \int_{y \in \mathbb{R}^d} \exp\left(\frac{v_0(y)}{h}\right) \exp\left(-\frac{|y-x|^2}{2t h}\right) dy \right) \end{aligned}$$

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Let

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Proposition – Value function [Lio82] The function

$$V(t,x) \coloneqq \int_{y \in \mathbb{R}^d}^{\oplus} v_0 \otimes \hat{\mu}_{t,x} = \sup_{y \in \mathbb{R}^d} \left[v_0(y) - \frac{|y-x|^2}{2t} \right]$$
(4)

is the unique viscosity solution of the Hamilton-Jacobi equation $\partial_t V - |\nabla V|^2 = 0$ such that $V(0, \cdot) = v_0$.

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is the unique viscosity solution of the Hamilton-Jacobi equation $\partial_t V - |\nabla V|^2 = 0$ such that $V(0, \cdot) = v_0$. The formula (4) is known as the Hopf-Lax formula.

The (max,+) semialgebra	Link with linear algebra	Finite-dim systems	The heat equation
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Going further			

• In the case of the heat equation, **Hopf-Cole** transform. But Lax-Hopf formula valid for a larger class of HJ equations of the type

$$\partial_t V(t,x) + H(x, \nabla V(t,x)) = 0,$$

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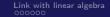
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- Maslov measures may be used to recast the Lax-Hopf semigroup as the conditional expectation of Maslov stochastic processes.
- Using the Hopf-Lax semigroup, Maslov defined weak solution by "duality", in the spirit of

$$\langle u, \varphi \rangle_{\oplus} = \langle u_0, S_t^* \varphi \rangle_{\oplus},$$

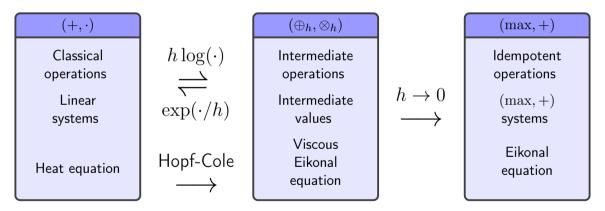
where S_t^* is a "dual" semigroup acting on test functions φ [KM97, Definition 3.1].

Conclusion



Finite-dim systems

The heat equation



Thank you!

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